

Examining the origin of dark matter mass at colliders

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As conventional dark matter scenarios have been probed extensively so far, the physics of a light dark matter charged under a new gauge group (dark gauge group) becomes one of new research avenues in many theoretical and experimental studies. We examine properties of a dark photon showering, the radiation process of light gauge bosons from energetic dark matter particles produced at the Large Hadron Collider (LHC). This showering process provides different signatures at the LHC depending on the property of dark matter under the dark gauge group. We show that the LHC experiment can identify the chirality of a dark matter, which leads to understanding the mass origin of particles in the dark sector.

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I. INTRODUCTION

The confirmation of dark matter (DM) existence will be the one of major milestones toward the physics beyond the standard model (SM) of the particle physics. Among various scenarios about dark matter, a weakly-interacting massive particle (WIMP) [1], often in the supersymmetry framework [2], has been extensively tested by various dark matter direct detection (DD) experiments [3,4] together with collider experiments including the LHC [5]. In near future, sensitivity of DD experiments will reach the point of detecting irreducible backgrounds from neutrino-nucleus coherent scattering [6].

In contrast to WIMP dark matter, searching for a light dark matter of sub-GeV mass scale with conventional DD experiments is very challenging due to low nuclear recoil energy $E_{\text{NR}} < \mathcal{O}(0.1)$ keV over experimental resolutions and noises [7]. There has been growing interest in the sub-GeV dark matter recently, and new experiments have been initiated and proposed [8–10]. They include direct searches of a relic dark matter particle [11] as well as beam experiments that produce

dark matter particles and detect their signals by using low-energy beam facilities [12,13]. As a contrast to DD experiments, the LHC has provided results of detecting a sub-GeV dark matter particle by utilizing initial state radiation (ISR) jet to tag events with dark matter particles as we have accumulated precise understanding in quantum chromodynamics (QCD) to suppress SM backgrounds [14].

Together with conventional dark matter experiments which are sensitive mostly to interactions between dark matter and SM particles, it would be interesting to consider phenomenological effects of dark matter scenarios if there exists an interaction among dark matter itself [15,16]. One of the natural methods to implement an interaction among dark matter particles is to introduce a dark gauge symmetry on dark matter [17–22]. In particular, a light dark gauge boson has been spotlighted in the intensity frontier research [23]. The combination of a light dark matter particle and a light gauge boson fits well, as a light gauge boson can provide a suitable annihilation channel for the dark matter particles demanded by the DM relic density constraint [24,25].

In this paper, we investigate generic collider signatures of a light dark matter particle, which is charged under a dark gauge group. The collider phenomenology of non-Abelian dark gauge group has been studied for a composite dark matter particle [26–28]. Here we focus on an Abelian dark gauge group, motivated by current efforts in dark photon searches [23]. Dark photon, as a dark gauge boson, can be produced through decay processes [29], or final state radiation [30,31]. Dark photon shower process can be triggered, once a “dark charged” dark matter particle is produced with a sufficient energy [32]. Our emphasis is to explain the difference in collider signatures, depending on

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the chirality of dark matter under a dark gauge group.¹ A chiral interaction is induced by a dark Higgs boson if it provides a mass to a dark photon through dark symmetry breaking and if dark matter becomes massive via Yukawa interaction. Thus we point out that recognizing patterns in dark photon showering can be a good probe to examine the mass origin of particles in a dark sector.

II. DARK SECTOR

Here we describe a minimal dark sector with a dark matter under a dark gauge group $U(1)_d$. If a dark sector contains a dark matter as a fundamental particle, the corresponding Lagrangian will be following;

$$\mathcal{L}_{\text{vector+scalar}} \ni -\frac{1}{4}F'_{\mu\nu}F'^{\mu\nu} + \frac{\varepsilon}{2}F_{\mu\nu}F'^{\mu\nu} + |D_\mu\Phi|^2, \quad (1)$$

$$\begin{aligned} \mathcal{L}_{\text{matter}} = & \bar{\chi}_L i\gamma^\mu D_\mu \chi_L + \bar{\chi}_R i\gamma^\mu D_\mu \chi_R + \bar{\psi}_L i\gamma^\mu D_\mu \psi_L \\ & + \bar{\psi}_R i\gamma^\mu D_\mu \psi_R - y_\chi \bar{\chi}_L \Phi^* \chi_R - y_\chi \bar{\chi}_R \Phi \chi_L \\ & - y_\psi \bar{\psi}_L \Phi \psi_R - y_\psi \bar{\psi}_R \Phi^* \psi_L, \end{aligned} \quad (2)$$

with $D_\mu \equiv \partial_\mu + ig'Q'A'_\mu$ where A'_μ is the quantum field of a dark photon γ_d , g' is the gauge coupling of the dark gauge symmetry, Q' is the dark $U(1)_d$ charge. $F_{\mu\nu}$ and $F'_{\mu\nu}$ are the field strength of the SM photon and dark photon respectively, and ε is the kinetic mixing parameter [34]. The SM particles are not charged under a dark gauge group. Here as we focus on a light dark photon where the mass of a dark photon is negligible compared to the mass of Z boson, the effect of SM electroweak symmetry breaking on a kinetic mixing becomes irrelevant to interactions between a dark photon and particles in the SM [35,36]. Φ is a dark Higgs which may break $U(1)_d$ depending on its charge under the dark gauge group. We introduce two pairs of chiral fermions χ as dark matter and ψ as a heavier particle in a dark sector for the anomaly cancellation, which is the model-dependent part. The Yukawa terms in Eq. (2) dictate the relations of the dark $U(1)_d$ charges;

$$Q'_{\chi_L} - Q'_{\chi_R} + Q'_\Phi = 0, \quad (3)$$

$$-Q'_{\psi_L} + Q'_{\psi_R} + Q'_\Phi = 0. \quad (4)$$

For the anomaly cancellation, we take $Q'_{\chi_L/\chi_R} = -Q'_{\psi_L/\psi_R}$, which allows a mixing between the χ and ψ . In our analysis, however, we will focus on the phenomenology of a dark matter χ . As pointed out in [33], nonzero charge Q'_Φ of a dark Higgs induces the chiral nature of dark matter. When we rewrite the interaction between a dark matter and a dark photon,

$$\mathcal{L}_{\text{matter}} \ni -g'Q'_VA'_\mu\bar{\chi}\gamma^\mu\chi - g'Q'_AA'_\mu\bar{\chi}\gamma^\mu\gamma_5\chi, \quad (5)$$

the axial Q'_A and vector coupling Q'_V are written as

$$Q'_A = \frac{1}{2}(Q'_{\chi_R} - Q'_{\chi_L}) = \frac{Q'_\Phi}{2} \quad (6)$$

$$Q'_V = \frac{1}{2}(Q'_{\chi_R} + Q'_{\chi_L}) = \frac{Q'_\Phi}{2} + Q'_{\chi_L}. \quad (7)$$

The chirality of a dark matter particle ($Q'_A \neq 0$) results in a significant difference in collider signatures at the LHC as we will show later.² Checking the chirality of a dark matter is directly related to understanding the mass generating mechanism for a dark matter.

There are several ways to address a small dark gauge boson mass even when a gauge coupling constant is not small [39–41]. In this paper we take a small vacuum expectation value v_S of the dark Higgs Φ given by $\Phi = \frac{1}{\sqrt{2}}(v_S + S + i\phi_S)$. The masses of the dark gauge boson and dark matter are given by the vacuum expectation value of the dark Higgs boson as $m_{\gamma_d} = g'Q'_\Phi v_S$ and $m_\chi = y_\chi v_S/\sqrt{2}$.

A. Dependence of dark photon showering on a mass mechanism in a dark sector

As an accelerated charged particle radiates corresponding gauge particles, energetic dark matter particles which are produced at a high energy collider will radiate off dark gauge bosons. The radiation pattern of a dark photon γ_d from an energetic dark matter χ , called the showering process, depends on the mechanism of the mass generation for a dark matter since dark matter couples differently with a dark photon as in Eq. (5). This showering process is characterized by a splitting function $P_{\chi \rightarrow \chi\gamma_d}$ which describes an emission process. In a collinear region, the differential probability of the splitting process $\chi \rightarrow \chi\gamma_d$ is

$$\frac{\alpha'}{2\pi} dx \frac{dt}{t} P_{\chi \rightarrow \chi\gamma_d}(x, t). \quad (8)$$

Here, $\alpha' = g'^2/4\pi$, t is the virtuality of incoming χ , and x is the energy fraction taken by outgoing χ . A detailed analysis of dark photon showers from vectorlike dark fermion model has been studied in [32]. As we focus on the phenomenology of an energetic dark matter production at colliders, we ignore terms suppressed by m_χ^2/t or $m_{\gamma_d}^2/t$. In this limit, the splitting kernel for vectorlike dark matter is given [42]

¹A systematic analysis of the dark mass origin and its impact on cosmology and astronomy is given in [33].

²This can be inferred from, for example, the distinct phenomenology depending on the chirality of SM fermions on dark gauge group [37,38] when corresponding dark gauge boson is very light.

$$P_{\chi \rightarrow \chi \gamma_d}(x, t) \simeq Q_V'^2 \frac{1+x^2}{1-x}. \quad (9)$$

This shower pattern is similar to the familiar QED shower which only includes contributions from transverse polarizations of a photon.

The longitudinal polarization vector of a dark photon will grow as $E_{\gamma_d}/m_{\gamma_d}$, with the energy of a dark photon E_{γ_d} . This artificial enhancement can be tamed by the Goldstone boson equivalence theorem (GBET) [43]. While the leading contribution of a longitudinal polarization in high energy limit is expressed by GBET which is proportional to m_χ/m_{γ_d} , the remaining part is suppressed by $m_{\gamma_d}/E_{\gamma_d}$ [44–46]. Thus this subleading part can be neglected in our study as we are interested in the phase space region of $E_{\gamma_d} \gg m_{\gamma_d}$. With GBET, we obtain a splitting kernel for the chiral fermions as following;

$$P_{\chi \rightarrow \chi \gamma_d}(x, t) \simeq (Q_V'^2 + Q_A'^2) \frac{1+x^2}{1-x} + 2Q_A'^2 \frac{m_\chi^2}{m_{\gamma_d}^2}. \quad (10)$$

The first term is from the transverse modes of a dark photon. The second term is from the longitudinal mode of a dark photon, which would be significant when a dark photon is very light compared to dark matter. Unlike the chiral dark matter, a vectorlike dark matter does not have an interaction with a Goldstone boson. Thus only transverse polarization of a dark photon is involved in a showering process as in Eq. (9). In the next section, we show the corresponding collider phenomenology by examining dark photon showering pattern.

III. PHENOMENOLOGY OF THE DARK SHOWERING AT THE LHC

The parameters that are directly related to a dark photon showering are $(\alpha', m_\chi, m_{\gamma_d})$ and (Q_A', Q_V') . In a case of chiral dark matter induced by nonzero $Q_\Phi' (= 1)$, we probe the case of $Q_{\chi_R}' = 0$ which maximizes effects of chirality with $Q_{\chi_L}' = -Q_\Phi' = -1$ as in Eq. (3). In terms of axial and vector coupling, we will have $(Q_A', Q_V') = (\frac{1}{2}, -\frac{1}{2})$ in this chiral case. For a vectorlike dark matter scenario, we consider $(Q_A', Q_V') = (0, 1)$ which comes from $(Q_{\chi_L}', Q_{\chi_R}') = (1, 1)$ as in Eq. (7). In this case, Q_Φ' becomes 0, which in turns decouples the origins of dark matter mass and dark photon mass.

The longitudinal component of a splitting kernel in Eq. (10) indicates that a large mass hierarchy between m_χ and m_{γ_d} will induce significant difference in the shower process from the chiral dark matter compared to the case of a vectorlike dark matter. But in order to keep the dark Yukawa coupling within a perturbative limit as $(y_\chi/\sqrt{2})^2 \lesssim 4\pi$, the mass spectrum of m_χ and m_{γ_d} follow a limit

TABLE I. Benchmark points we have chosen. They obey the perturbative limit of $\alpha' \frac{m_\chi^2}{m_{\gamma_d}^2} \lesssim 1$.

Benchmark Points (BP)	A	B	C
α'	0.3	0.15	0.075
m_χ (GeV)	0.7	1.0	1.4
m_{γ_d} (GeV)	0.4		

$$\alpha' \frac{m_\chi^2}{m_{\gamma_d}^2} \lesssim 1. \quad (11)$$

To avoid constraints from current dark matter and dark photon search experiments [3,4,7,8,10–13] we take quite light benchmark points as in Table I. In our benchmark points of $m_{\gamma_d} = 0.4$ GeV, a viable kinematic mixing parameter ϵ can provide prompt decays of dark photons to the SM particles. More specifically, for our benchmark points, ϵ^2 smaller than 10^{-7} is still allowed by current constraints [9]. If ϵ is too small, SM particles from a dark photon decays would leave displaced vertices for us to enhance the search power at the LHC [47]. With $\epsilon^2 > 10^{-10}$, the impact parameter of particles from dark photon is smaller than 1 mm, as particles from dark photon decays can be treated as prompt [48]. In this case, a dark photon mostly decays into a pair of light leptons where these non-conventional signatures are easy to tag over the QCD backgrounds. The corresponding branching ratios are $\text{BR}(\gamma_d \rightarrow \mu^+ \mu^-) \simeq 0.45$, $\text{BR}(\gamma_d \rightarrow e^+ e^-) \simeq 0.45$, and $\text{BR}(\gamma_d \rightarrow \pi^+ \pi^-) \simeq 0.10$. From now on, we fix the mass of γ_d to 0.4 GeV.

Particles from dark photon showering processes would be tagged at collider detectors when they can leave certain level of energy deposits. Since a dark photon showering process as a final state radiation of dark matter is insensitive to a production process, we consider a TeV-scale mediator Z' in producing dark matter particles at a collider in Fig. 1. As null results of the LHC push the possible mass range of a mediator to be heavy, our set up in the framework of “hidden valley” [49] is empirically supported.³ For a phenomenological study in a hadron collider, we take minimal interactions between the standard model sector and a dark sector when the mass of Z' is within the coverage of the high luminosity (HL) of the LHC

$$\mathcal{L} \ni -g_q Z'_\mu \bar{q} \gamma^\mu q + g_\chi Z'_\mu \bar{\chi} \gamma^\mu \chi. \quad (12)$$

here q denotes a SM quark. If the energy of a hadron collider is not enough to produce a on-shell mediator, dark matter productions will be described by an effective operator with a mediator being integrated out [50].

³Considering a very heavy Z' can be introduced as a mechanism to make γ_d very light through mass a matrix diagonalization [41].

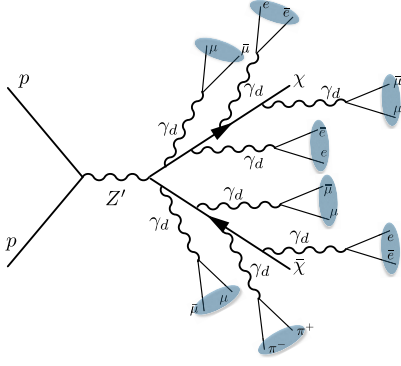


FIG. 1. A High energy collider produces dark matter particles, which get enough energy recoils from a decay of a heavy resonance to shower dark photons.

$$\mathcal{L} \ni \frac{1}{\Lambda^2} (\bar{q} \gamma^\mu q) (\bar{\chi} \gamma_\mu \chi). \quad (13)$$

In this case, a phase space of events with boosted dark matter is different from our current study as an initial state radiation jet will be the source of producing boosted dark matters. Here we focus on a situation where HL-LHC can reach the mass range of a mediator with $M_{Z'} = 1.5$ TeV. Model parameters of coupling constants are chosen to be compatible with current LHC searches of dijet and prompt lepton-jet searches as we show later. We perform Monte Carlo studies with FEYNRULES2.0 [51] to implement dark matter models, MADGRAPH_AMC@NLO [52] and PYTHIA8 [53]. In simulating dark photon showering processes, we modify a hidden valley model [49,54] implemented in PYTHIA8 [55,56] to add longitudinal term in Eq. (10).

To examine a difference in signatures at the LHC from distinct showering patterns, we check how many changes occur in the number of produced dark photons at the LHC depending on the chirality of dark matter. In the left column of Fig. 2, we plot histograms of the number of showered dark photons per event. As we observe, with increasing the mass of dark matter, the number of showered dark photons is reduced in a vectorlike dark matter case as a dark gauge coupling α' is decreasing as in our benchmark points in Table I. But in a chiral dark matter scenario, the number of the dark photon is almost unchanged due to the enhancement from the GBET with a large Yukawa coupling in the second term of Eq. (10).

Once, a dark photon is produced from a dark showering process, it decays to SM particles through a kinetic mixing ϵ in Eq. (1). To quantify those signatures, we use the scalar sum of the muon and electron transverse momentum p_T ;

$$H_{T(\ell)} = \sum_{i=\mu^\pm, e^\pm} |p_{Ti}|. \quad (14)$$

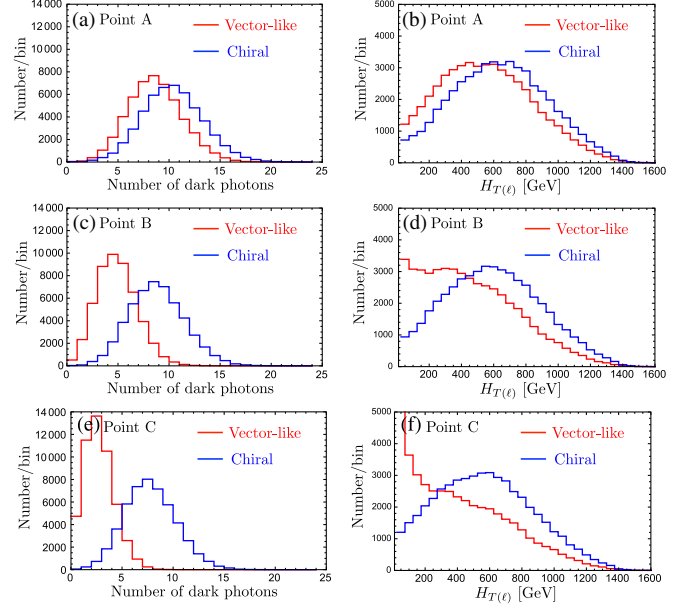


FIG. 2. To show different dark photon shower pattern, we demonstrate the number of showered dark photons and corresponding $H_{T(\ell)}$ distributions with parton level Monte Carlo simulations. The left column is for the comparison of the dark photon number distribution of our benchmark points between a vectorlike dark matter and chiral dark matter. Red line and blue line correspond to vectorlike model and chiral model respectively. The right column is for the comparison of the $H_{T(\ell)}$ distribution of our benchmark points with the same color code as cases in the left column.

In the right column of Fig. 2 we observe that a dark photon shower in the vectorlike dark matter case becomes weak and most of leptons become soft as a coupling constant α' gets smaller. In a chiral dark matter case, showering processes with a longitudinal mode of a dark photon, which is independent on α' , becomes dominant through an enhancement from a Yukawa interaction as dark matter becomes massive. This can be shown in the following limit

$$\lim_{\alpha' \ll 1} \alpha' P_{\chi \rightarrow \chi \gamma_d} \simeq \alpha' \cdot 2Q_A'^2 \frac{m_\chi^2}{m_{\gamma_d}^2} \propto \frac{m_\chi^2}{v_S^2} \cdot \frac{m_\chi^2}{m_{\gamma_d}^2} \sim y_\chi^2. \quad (15)$$

Due to GBET, the energy spectrum of leptons from a longitudinal mode is larger compared to the case of leptons from a transverse mode of dark photons.

To consider various factors including isolation of reconstructed objects and smearing effects on energy deposits, we use Delphes 3 [57] as a fast detector simulation with the ATLAS parameter setting. We adopt the concept of lepton-jet (LJ) [47,58] to cluster collimated muons from a light dark photon. As a light dark photon decays mostly the pair of leptons, we only consider a muon-jet (LJ_μ) as a candidate to suppress QCD backgrounds as we can identify individual muons in LJ_μ by utilizing various subdetectors to

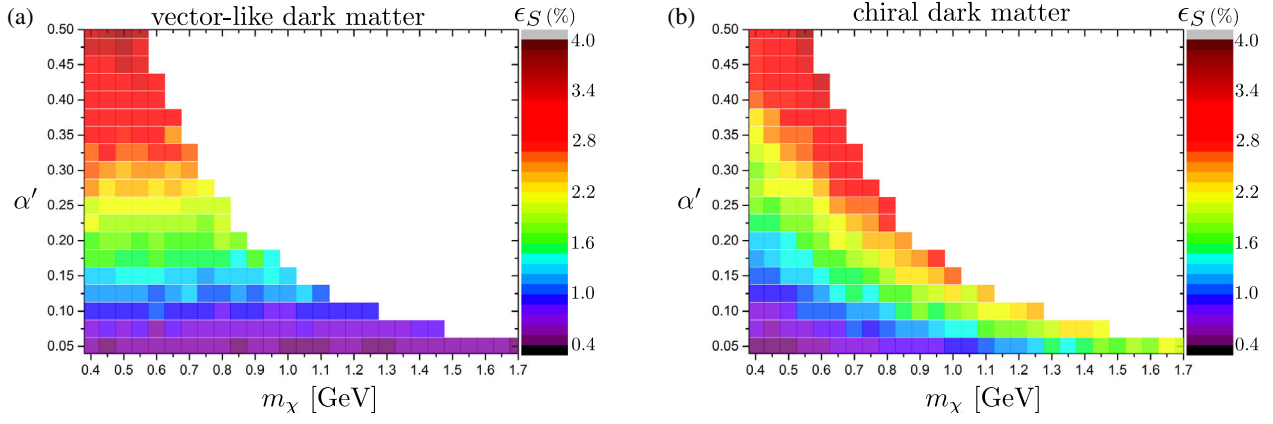


FIG. 3. We show the signal efficiency ϵ_S of tagging two muon-jets as in a search of prompt decaying light bosons of ATLAS [65].

perform χ^2 fitting for muon tracks. Considering electron-jet (LJ_e) for a signal object would nontrivial as bremsstrahlung processes of electrons force us to merge several crystals in electronic calorimeters, making it difficult to isolate each electrons in collimated situations and requiring additional properties to reduce QCD backgrounds [47,59]. Thus we require at least two muon-jets (LJ_μ) to tag signals over backgrounds. After we require an isolation criteria for a muon-jet, we choose the mass of muon-jet (LJ_μ) within the range between 0.3 GeV and 0.5 GeV to consider imperfect resolution of detectors [60].

A. Constraints from the LHC searches

In this section, we present constraints on our benchmark scenarios provide by the LHC searches. As we introduce a heavy leptophobic Z' as the mediator of a simplified model for a collider phenomenology, there is a limit from a heavy resonance search in a dijet signature [61,62]. The difference between two LHC searches (ATLAS and CMS) for a heavy Z' is that they have different coupling structure between a mediator and each sector (standard model sector and dark sector). ATLAS search assumes an axial-vector coupling between Z' and quarks, dark matter, motivated by negative results from dark matter direct searches [63], while CMS analysis takes a vector coupling of a model of $U(1)_B$ associated with gauged baryon number [64]. Different coupling structure can affect cut-efficiency as an angular distribution depends on it. But as the difference is proportional to a ratio, $m_d^2/M_{Z'}^2$, where m_d is the mass of particles from Z' decaying, we can safely combine results from ATLAS and CMS together in recasting these analyses to our case with $M_{Z'} \gg m_q, m_\chi$.

To apply constraints from dijet resonance searches, we consider the invisible branch ratio of Z' . With $M_{Z'} = 1.5$ TeV, the most stringent upper limit on a coupling g_q between Z' and the standard model quarks q is from the ATLAS as $g_q \lesssim 0.07$. This upper limit comes

from narrow-width case where Z' decays only to light quarks. A constraint becomes milder with increasing Z' width due to losing sensitivity in wider dijet mass window [61]. Thus our estimation with a narrow-width Z' is conservative. Sizable invisible decay partial width from a coupling g_χ between dark matter and Z' , the limit on (g_q, g_χ) can be imposed by

$$g_q^2 \times \frac{N_c \cdot N_{\text{lf}} \cdot g_q^2}{N_c \cdot N_f \cdot g_q^2 + g_\chi^2} \lesssim 0.07^2 \times \frac{N_c \cdot N_{\text{lf}} \cdot g_q^2}{N_c \cdot N_f \cdot g_q^2}, \quad (16)$$

with $N_c = 3$ is the color factor of the standard model and $N_{\text{lf}} = 5$ is the number of light flavor quarks considered as the final state jets in the LHC dijet searches. Here $N_f = 5 + \sqrt{1 - 4m_t^2/M_{Z'}^2} \simeq 5.97$ is the effective number of quark flavors contributing to the width of Z' [62]. As we see in the left side of Eq. (16), the constraint from dijet searches become weaker as Z' has non-negligible invisible decay proportional to g_χ^2 . As the (partial) decay width of Z' is not sensitive to the mass of sub-GeV dark matter, a constraint from dijet searches will not depend on m_χ .

We also consider constraints from prompt lepton-jet analysis [65]. Out of various combinations of lepton-jets as a signal channel, we consider two muon-jets signal as this channel has the highest tagging efficiency and lowest backgrounds [65,66]. In recasting a prompt lepton-jet analysis, we take two extreme cases where (a) the parameter space allows the maximum dark photon showering activities, and (b) we have least dark photon showering activities within the perturbative limit of $\alpha' m_\chi^2/m_{\gamma_d}^2 \lesssim 1$ as in Table II. These two limiting cases provide the upper and lower bounds of the signal efficiency ϵ_S in tagging two isolated muon-jets as in Fig. 3. Thus we can see the maximally and minimally allowed range of g_χ from the prompt lepton-jet analyses in our parameter space of (α', m_χ) by considering these two extreme cases.

TABLE II. Benchmark points to consider ATLAS prompt lepton-jet (LJ) analysis conservatively. For (a) strong coupling limit, we take a vectorlike dark matter with $(m_\chi, \alpha') = (0.5 \text{ GeV}, 0.5)$ and for (b) weak coupling limit, we choose a chiral dark matter with $(m_\chi, \alpha') = (0.45 \text{ GeV}, 0.05)$. Here ϵ_S is a signal efficiency in tagging at least two isolated muon-jets (LJ_μ) from the ATLAS [65].

Benchmark case	α'	$m_\chi(\text{GeV})$	Model	$\epsilon_S(\%)$
(a) Strong coupling limit	0.50	0.5	Vector-like	3.68
(b) Weak coupling limit	0.05	0.45	Chiral	0.41

B. Analysis at the High Luminosity (HL) LHC

In this section, we provide a viable parameter space to have enough statistics in discriminating two different hypothesis on the property of dark matter. For the HL-LHC analysis, we consider a muon-jet as following. First of all, a candidate muon in the muon-jet needs to have $p_T > 5 \text{ GeV}$ within $|\eta| < 2.5$. Then we use the Cambridge/Aachen jet algorithm [67] with FASTJET [68] to cluster a muon-jet with $\Delta R < 0.1$. An isolation variable for a muon-jet is defined as:

$$\rho = \frac{\sum_i E_{T,i}}{p_{T,\mu J}}, \quad (17)$$

with i running over all the E_T deposit in the calorimeter, without the candidate muons, near the muon-jet with $\Delta R < 0.3$, and the denominator is the p_T of the muon-jet. To suppress the QCD faking rate, we require the isolation criteria as $\rho < 0.3$. After that, the major backgrounds in our case where we require at least two prompt muon-jets would be low-mass mesons including ρ , ω and ϕ which decay into muon pair [65,66].⁴ To reduce these backgrounds, we require the mass of reconstructed muon-jet within (0.3,0.5) in the unit of [GeV]. This requirement reduces backgrounds from low-mass mesons, leaving the di-photon process ($pp \rightarrow \gamma^*$, $\gamma^* \rightarrow 2\text{LJ}_\mu$) as the major background. To suppress contribution from a di-photon process, we require $H_{T(\ell)} \geq 100 \text{ GeV}$. As we see in Fig. 4, the most of events in background locate at low mass region of $H_{T(\ell)}$. By requiring a cut on $H_{T(\ell)}$, we have only 3.07 events from backgrounds at 14 TeV HL-LHC. We show the parameter space of (g_q, g_χ) which has not been excluded by current LHC analyses and will have enough statistics to distinguish vectorlike and chiral dark matter at the HL-LHC of $\mathcal{L} = 3 \text{ ab}^{-1}$ in Fig. 5. Constraints from dijet analysis do not depend on α' as the analysis focuses on dijet final states. It becomes weak as g_χ becomes large

⁴In the parameter space of nonpromptly decaying light particles, $b\bar{b}$ production becomes a major background as b -quarks decay into pairs of muons via double semileptonic decays [69].

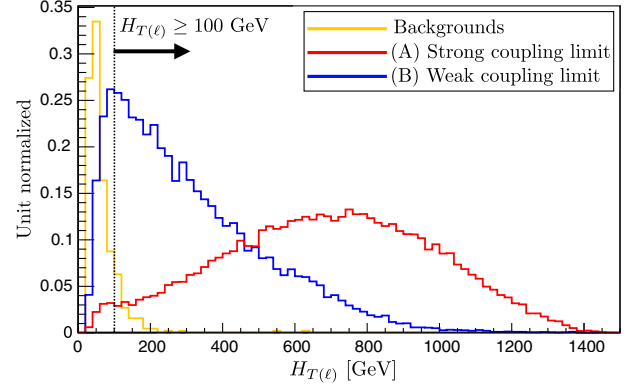


FIG. 4. Distributions of $H_{T(\ell)}$ after requiring isolation and mass window on reconstructed muon-jets. The dominant background is from the diphoton production where off-shell photons γ^* decays into muon pair. The background distribution is shown with a yellow line. we require a cut on $H_{T(\ell)}$ to suppress backgrounds accordingly for the HL-LHC study.

enough to induce a significant invisible decay partial width of Z' . The number of events N_{sig} depends on α' as the signal tagging efficiency depends on the amount of dark photon showering which is proportional to α' in vectorlike dark matter case. As it is shown in Fig. 5, there would be viable parameter space where we can see the relation between the mass of dark matter and the mass of dark photon by examining the dark photon shower pattern. Finally, we describe how one can understand the nature of a dark sector once we observe dark matter signatures at the LHC. For a collider observable to identify the nature of dark matter, we compare $H_{T(\ell)}$ variable in Eq. (14) distribution with muons of $p_T \geq 5 \text{ GeV}$ and electrons of $p_T \geq 10 \text{ GeV}$ by considering momentums from tracks and energy deposits in subdetectors. To suppress backgrounds which populate the low $H_{T(\ell)}$ region, we consider only high mass region of $H_{T(\ell)} \geq 100 \text{ GeV}$. In our simulation, we have $\mathcal{O}(1)$ events from backgrounds at HL-LHC. To deal with a finite luminosity of the LHC, we generated 200 reconstructed signal events after cuts. We perform 100 pseudo-experiments to reduce statistical uncertainties coming from finite statistics of simulations. We calculate a binned- χ^2 of $H_{T(\ell)}$ distributions from a vectorlike case and chiral case with including corruptions from backgrounds as we described above. Figure 6 shows our results from χ^2 comparison. The LHC can tell the origin of the mass of a dark matter particle by discriminating a vectorlike and chiral dark matter models more than 2σ significance level, when mass ratio m_χ/m_{γ_d} is large enough. While with a significant events number, as in Fig. 6 (c), for parameter region of large α' and small m_χ/m_{γ_d} , the vectorlike and chiral dark matter models can be also distinguished by 2σ significance level because of the weaker transverse modes shower in the chiral dark matter case.

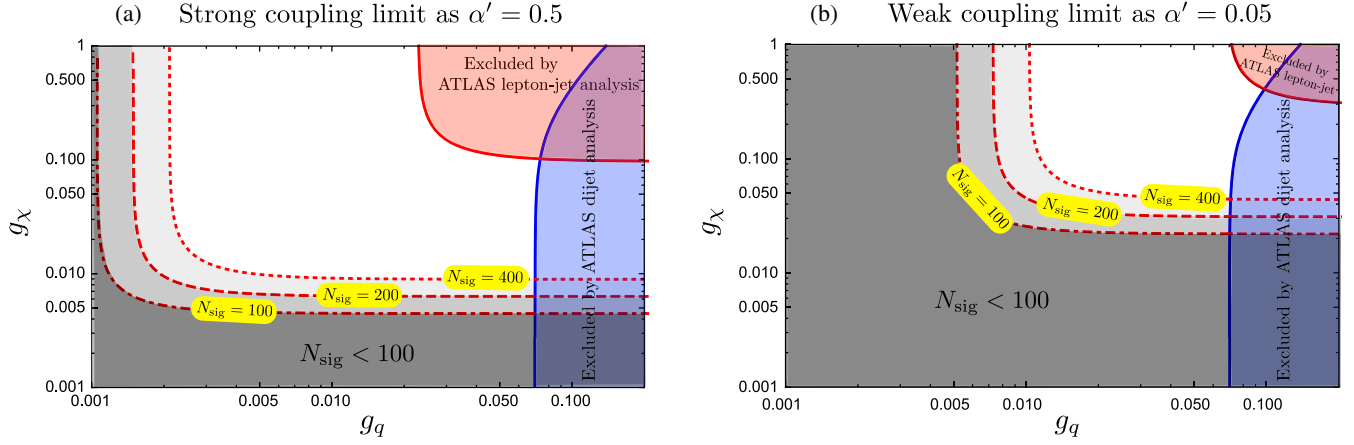


FIG. 5. We show possible parameter space of (g_q, g_χ) with moderate events after analysis cuts at the HL-LHC of $\mathcal{L} = 3 \text{ ab}^{-1}$ to discriminate chiral dark matter from vectorlike dark matter. The blue shaded region is excluded by the ATLAS dijet searches. The red region is excluded by the ATLAS prompt lepton-jet search. The grey shaded regions are divided by the number of signal events N_{sig} less than 100, 200, and 400 after analysis cuts.

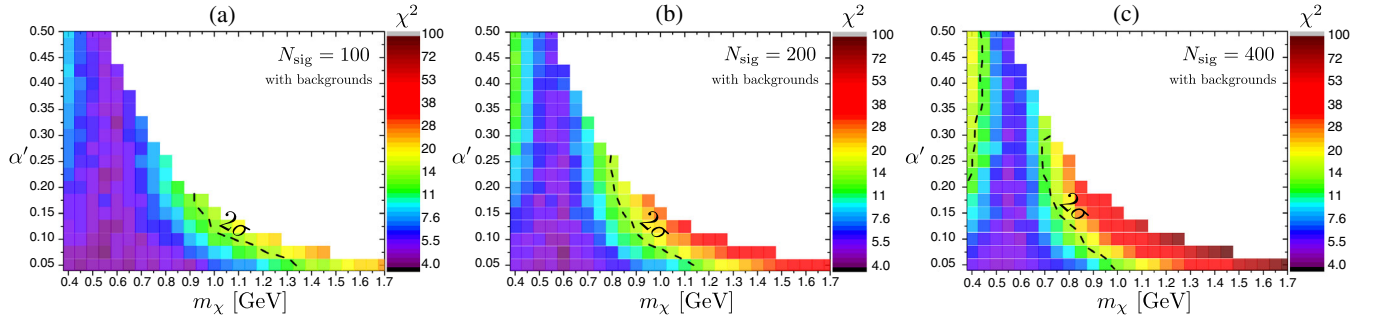


FIG. 6. χ^2 distance of $H_{T(\ell)}$ distributions between a vectorlike dark matter scenario and chiral dark matter case with 100, 200, and 400 signal events with backgrounds after cuts.

IV. DISCUSSIONS AND OUTLOOK

The standard model had been developed by constructing gauge structures. At the LHC, we checked its validity from the discovery of the Higgs particle which provides a mass to SM particles. Similarly, identifying the gauge structure of a dark sector together with understanding the mass origin of dark particles would be the first step toward expanding the physics of dark matter once we observe a dark matter signature. With the performance of the LHC which is a complementary tool to probe a light dark matter, we study the feature of collider signatures from a dark photon showering depending on the property of dark matter under a dark gauge group. With numerical simulations of 14 TeV

LHC, we show that we can identify the nature of dark matter with $\mathcal{O}(100)$ signal events at high energy colliders.

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- [1] B. W. Lee and S. Weinberg, Cosmological Lower Bound on Heavy Neutrino Masses, *Phys. Rev. Lett.* **39**, 165 (1977).
- [2] G. Jungman, M. Kamionkowski, and K. Griest, Supersymmetric dark matter, *Phys. Rep.* **267**, 195 (1996); G. Bertone, D. Hooper, and J. Silk, Particle dark matter: Evidence, candidates and constraints, *Phys. Rep.* **405**, 279 (2005).
- [3] D. S. Akerib *et al.* (LUX Collaboration), Results from a Search for Dark Matter in the Complete LUX Exposure, *Phys. Rev. Lett.* **118**, 021303 (2017); Limits on Spin-Dependent WIMP-Nucleon Cross Section Obtained from the Complete LUX Exposure, *Phys. Rev. Lett.* **118**, 251302 (2017); X. Cui *et al.* (PandaX-II Collaboration), Dark Matter Results From 54-Ton-Day Exposure of PandaX-II Experiment, *Phys. Rev. Lett.* **119**, 181302 (2017); C. Fu *et al.* (PandaX-II Collaboration), Spin-Dependent Weakly-Interacting-Massive-Particle?Nucleon Cross Section Limits from First Data of PandaX-II Experiment, *Phys. Rev. Lett.* **118**, 071301 (2017); Erratum, **120**, 049902 (2018); E. Aprile *et al.* (XENON Collaboration), First Dark Matter Search Results from the XENON1T Experiment, *Phys. Rev. Lett.* **119**, 181301 (2017); E. Aprile *et al.* (XENON100 Collaboration), XENON100 dark matter results from a combination of 477 live days, *Phys. Rev. D* **94**, 122001 (2016).
- [4] R. Agnese *et al.* (SuperCDMS Collaboration), New Results from the Search for Low-Mass Weakly Interacting Massive Particles with the CDMS Low Ionization Threshold Experiment, *Phys. Rev. Lett.* **116**, 071301 (2016); D. S. Akerib *et al.* (LZ Collaboration), LUX-ZEPLIN (LZ) conceptual design report, [arXiv:1509.02910](#); B. J. Mount *et al.*, LUX-ZEPLIN (LZ) technical design report, [arXiv:1703.09144](#); J. Aalbers *et al.* (DARWIN Collaboration), DARWIN: Towards the ultimate dark matter detector, [arXiv:1606.07001](#).
- [5] V. Khachatryan *et al.* (CMS Collaboration), Search for dark matter particles in proton-proton collisions at $\sqrt{s} = 8$ TeV using the razor variables, *J. High Energy Phys.* **12** (2016) 088; G. Aad *et al.* (ATLAS Collaboration), Search for Dark Matter in Events with a Hadronically Decaying W or Z Boson and Missing Transverse Momentum in pp Collisions at $\sqrt{s} = 8$ TeV with the ATLAS Detector, *Phys. Rev. Lett.* **112**, 041802 (2014); D. Abercrombie *et al.*, Dark matter benchmark models for early LHC Run-2 searches: Report of the ATLAS/CMS dark matter forum, [arXiv:1507.00966](#); M. Aaboud *et al.* (ATLAS Collaboration), Search for dark matter at $\sqrt{s} = 13$ TeV in final states containing an energetic photon and large missing transverse momentum with the ATLAS detector, *Eur. Phys. J. C* **77**, 393 (2017); Search for Dark Matter Produced in Association with a Higgs Boson Decaying to $b\bar{b}$ using 36 fb^{-1} of pp collisions at $\sqrt{s} = 13$ TeV with the ATLAS Detector, *Phys. Rev. Lett.* **119**, 181804 (2017).
- [6] J. Billard, E. Figueroa-Feliciano, and L. Strigari, Implication of neutrino backgrounds on the reach of next generation dark matter direct detection experiments, *Phys. Rev. D* **89**, 023524 (2014).
- [7] G. Angloher *et al.* (CRESST Collaboration), Results on light dark matter particles with a low-threshold CRESST-II detector, *Eur. Phys. J. C* **76**, 25 (2016); R. Agnese *et al.* (SuperCDMS Collaboration), New Results from the Search for Low-Mass Weakly Interacting Massive Particles with the CDMS Low Ionization Threshold Experiment, *Phys. Rev. Lett.* **116**, 071301 (2016); Projected sensitivity of the SuperCDMS SNOLAB experiment, *Phys. Rev. D* **95**, 082002 (2017); A. Aguilar-Arevalo *et al.* (DAMIC Collaboration), Search for low-mass WIMPs in a 0.6 kg day exposure of the DAMIC experiment at SNOLAB, *Phys. Rev. D* **94**, 082006 (2016).
- [8] M. Battaglieri *et al.* (BDX Collaboration), Dark matter search in a Beam-Dump eXperiment (BDX) at Jefferson Lab: an update on PR12-16-001, [arXiv:1712.01518](#).
- [9] J. Alexander *et al.*, Dark Sectors 2016 Workshop: Community report, [arXiv:1608.08632](#); M. Battaglieri *et al.*, US cosmic visions: New ideas in dark matter 2017: Community report, [arXiv:1707.04591](#).
- [10] P. deNiverville, C. Y. Chen, M. Pospelov, and A. Ritz, Light dark matter in neutrino beams: Production modelling and scattering signatures at MiniBooNE, T2K and SHiP, *Phys. Rev. D* **95**, 035006 (2017).
- [11] R. Essig, A. Manalaysay, J. Mardon, P. Sorensen, and T. Volansky, First Direct Detection Limits on sub-GeV Dark Matter from XENON10, *Phys. Rev. Lett.* **109**, 021301 (2012); R. Essig, T. Volansky, and T. T. Yu, New constraints and prospects for sub-GeV dark matter scattering off electrons in xenon, *Phys. Rev. D* **96**, 043017 (2017).
- [12] A. A. Aguilar-Arevalo *et al.* (MiniBooNE DM Collaboration), Dark matter search in nucleon, pion, and electron channels from a proton beam dump with MiniBooNE, [arXiv:1807.06137](#).
- [13] D. Banerjee *et al.* (NA64 Collaboration), Search for Invisible Decays of Sub-GeV Dark Photons in Missing-Energy Events at the CERN SPS, *Phys. Rev. Lett.* **118**, 011802 (2017); Search for vector mediator of Dark Matter production in invisible decay mode, *Phys. Rev. D* **97**, 072002 (2018).
- [14] A. M. Sirunyan *et al.* (CMS Collaboration), Search for new physics in final states with an energetic jet or a hadronically decaying W or Z boson and transverse momentum imbalance at $\sqrt{s} = 13$ TeV, *Phys. Rev. D* **97**, 092005 (2018); M. Aaboud *et al.* (ATLAS Collaboration), Search for dark matter and other new phenomena in events with an energetic jet and large missing transverse momentum using the ATLAS detector, *J. High Energy Phys.* **01** (2018) 126.
- [15] B. A. Gradwohl and J. A. Frieman, Dark matter, long range forces, and large scale structure, *Astrophys. J.* **398**, 407 (1992).
- [16] D. N. Spergel and P. J. Steinhardt, Observational Evidence for Self-Interacting Cold Dark Matter, *Phys. Rev. Lett.* **84**, 3760 (2000).
- [17] D. Hooper and K. M. Zurek, A natural supersymmetric model with MeV dark matter, *Phys. Rev. D* **77**, 087302 (2008).
- [18] J. L. Feng and J. Kumar, The WIMPless Miracle: Dark-Matter Particles without Weak-Scale Masses or Weak Interactions, *Phys. Rev. Lett.* **101**, 231301 (2008).
- [19] L. Ackerman, M. R. Buckley, S. M. Carroll, and M. Kamionkowski, Dark matter and dark radiation, *Phys. Rev. D* **79**, 023519 (2009).
- [20] J. L. Feng, M. Kaplinghat, H. Tu, and H. B. Yu, Hidden charged dark matter, *J. Cosmol. Astropart. Phys.* **07** (2009) 004.
- [21] Y. Hochberg, E. Kuflik, H. Murayama, T. Volansky, and J. G. Wacker, Model for Thermal Relic Dark Matter

- of Strongly Interacting Massive Particles, *Phys. Rev. Lett.* **115**, 021301 (2015).
- [22] J. Kopp, J. Liu, T. R. Slatyer, X. P. Wang, and W. Xue, Impeded dark matter, *J. High Energy Phys.* **12** (2016) 033.
- [23] R. Essig *et al.*, Working group report: New light weakly coupled particles, [arXiv:1311.0029](https://arxiv.org/abs/1311.0029); Search for the dark photon in π^0 decays, *Phys. Lett. B* **746**, 178 (2015); S. Alekhin *et al.*, A facility to search for hidden particles at the CERN SPS: The SHiP physics case, *Rep. Prog. Phys.* **79**, 124201 (2016); P. Ilten, Y. Soreq, J. Thaler, M. Williams, and W. Xue, Proposed Inclusive Dark Photon Search at LHCb, *Phys. Rev. Lett.* **116**, 251803 (2016).
- [24] C. Boehm and P. Fayet, Scalar dark matter candidates, *Nucl. Phys.* **B683**, 219 (2004).
- [25] P. Fayet, Light spin 1/2 or spin 0 dark matter particles, *Phys. Rev. D* **70**, 023514 (2004).
- [26] P. Schwaller, D. Stolarski, and A. Weiler, Emerging jets, *J. High Energy Phys.* **05** (2015) 059.
- [27] T. Cohen, M. Lisanti, and H. K. Lou, Semivisible Jets: Dark Matter Undercover at the LHC, *Phys. Rev. Lett.* **115**, 171804 (2015).
- [28] C. Englert, K. Nordstrom, and M. Spannowsky, Towards resolving strongly-interacting dark sectors at colliders, *Phys. Rev. D* **94**, 055028 (2016).
- [29] K. Kong, H. S. Lee, and M. Park, Dark decay of the top quark, *Phys. Rev. D* **89**, 074007 (2014).
- [30] C. Cheung, J. T. Ruderman, L. T. Wang, and I. Yavin, Lepton jets in (supersymmetric) electroweak processes, *J. High Energy Phys.* **04** (2010) 116; A. Falkowski, J. T. Ruderman, T. Volansky, and J. Zupan, Hidden Higgs decaying to lepton jets, *J. High Energy Phys.* **05** (2010) 077; Y. F. Chan, M. Low, D. E. Morrissey, and A. P. Spray, LHC signatures of a minimal supersymmetric hidden valley, *J. High Energy Phys.* **05** (2012) 155; D. Curtin, R. Essig, S. Gori, and J. Shelton, Illuminating dark photons with high-energy colliders, *J. High Energy Phys.* **02** (2015) 157.
- [31] Y. Bai, J. Bourbeau, and T. Lin, Dark matter searches with a mono- Z^0 jet, *J. High Energy Phys.* **06** (2015) 205; A. Gupta, R. Primulando, and P. Saraswat, A new probe of dark sector dynamics at the LHC, *J. High Energy Phys.* **09** (2015) 079.
- [32] M. Buschmann, J. Kopp, J. Liu, and P. A. N. Machado, Lepton jets from radiating dark matter, *J. High Energy Phys.* **07** (2015) 045.
- [33] N. F. Bell, Y. Cai, and R. K. Leane, Impact of mass generation for spin-1 mediator simplified models, *J. Cosmol. Astropart. Phys.* **01** (2017) 039.
- [34] B. Holdom, Two $U(1)$'s and epsilon charge shifts, *Phys. Lett.* **166B**, 196 (1986).
- [35] E. J. Chun, J. C. Park, and S. Scopel, Dark matter and a new gauge boson through kinetic mixing, *J. High Energy Phys.* **02** (2011) 100.
- [36] H. S. Lee and S. Yun, Mini force: The $(B - L) + xY$ gauge interaction with a light mediator, *Phys. Rev. D* **93**, 115028 (2016).
- [37] H. Davoudiasl, H. S. Lee, and W. J. Marciano, 'Dark' Z implications for parity violation, rare meson decays, and Higgs physics, *Phys. Rev. D* **85**, 115019 (2012).
- [38] H. Davoudiasl, H. S. Lee, and W. J. Marciano, Muon Anomaly and Dark Parity Violation, *Phys. Rev. Lett.* **109**, 031802 (2012).
- [39] N. Arkani-Hamed and N. Weiner, LHC signals for a superunified theory of dark matter, *J. High Energy Phys.* **12** (2008) 104.
- [40] C. Cheung, J. T. Ruderman, L. T. Wang, and I. Yavin, Kinetic mixing as the origin of light dark scales, *Phys. Rev. D* **80**, 035008 (2009).
- [41] H. S. Lee and M. S. Seo, Gauge see-saw: A mechanism for a light gauge boson, *Phys. Lett. B* **767**, 69 (2017).
- [42] S. Catani, S. Dittmaier, M. H. Seymour, and Z. Trocsanyi, The dipole formalism for next-to-leading order QCD calculations with massive partons, *Nucl. Phys.* **B627**, 189 (2002).
- [43] J. M. Cornwall, D. N. Levin, and G. Tiktopoulos, Derivation of gauge invariance from high-energy unitarity bounds on the S matrix, *Phys. Rev. D* **10**, 1145 (1974); C. E. Vayonakis, Born helicity amplitudes and cross-sections in non-Abelian gauge theories, *Lett. Nuovo Cimento Soc. Ital. Fis.* **17**, 383 (1976); B. W. Lee, C. Quigg, and H. B. Thacker, Weak interactions at very high energies: The role of the Higgs-boson mass, *Phys. Rev. D* **16**, 1519 (1977); M. S. Chanowitz and M. K. Gaillard, The TeV physics of strongly interacting W 's and Z 's, *Nucl. Phys.* **B261**, 379 (1985); J. Bagger and C. Schmidt, Equivalence theorem redux, *Phys. Rev. D* **41**, 264 (1990); H. Veltman, The equivalence theorem, *Phys. Rev. D* **41**, 2294 (1990); W. Beenakker and A. Werthenbach, Electroweak two-loop Sudakov logarithms for on-shell fermions and bosons, *Nucl. Phys.* **B630**, 3 (2002);
- [44] A. Wulzer, An equivalent gauge and the equivalence theorem, *Nucl. Phys.* **B885**, 97 (2014).
- [45] P. P. Srivastava and S. J. Brodsky, A unitary and renormalizable theory of the standard model in ghost free light cone gauge, *Phys. Rev. D* **66**, 045019 (2002).
- [46] J. Chen, T. Han, and B. Tweedie, Electroweak splitting functions and high energy showering, *J. High Energy Phys.* **11** (2017) 093.
- [47] G. Aad *et al.* (ATLAS Collaboration), Search for long-lived neutral particles decaying into lepton jets in proton-proton collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector, *J. High Energy Phys.* **11** (2014) 088; ATLAS Collaboration, Search for long-lived neutral particles decaying into displaced lepton jets in proton-proton collisions at $\sqrt{s} = 13$ TeV with the ATLAS detector, Report No. ATLAS-CONF-2016-042.
- [48] G. Aad *et al.* (ATLAS Collaboration), Electron reconstruction and identification efficiency measurements with the ATLAS detector using the 2011 LHC proton-proton collision data, *Eur. Phys. J. C* **74**, 2941 (2014).
- [49] M. J. Strassler and K. M. Zurek, Echoes of a hidden valley at hadron colliders, *Phys. Lett. B* **651**, 374 (2007).
- [50] J. Goodman, M. Ibe, A. Rajaraman, W. Shepherd, T. M. P. Tait, and H. B. Yu, Constraints on dark matter from colliders, *Phys. Rev. D* **82**, 116010 (2010).
- [51] A. Alloul, N. D. Christensen, C. Degrande, C. Duhr, and B. Fuks, FeynRules 2.0—A complete toolbox for tree-level phenomenology, *Comput. Phys. Commun.* **185**, 2250 (2014).
- [52] J. Alwall, R. Frederix, S. Frixione, V. Hirschi, F. Maltoni, O. Mattelaer, H.-S. Shao, T. Stelzer, P. Torrielli, and M. Zaro, The automated computation of tree-level and next-to-leading order differential cross sections, and their matching to parton shower simulations, *J. High Energy Phys.* **07** (2014) 079.

- [53] T. Sjöstrand, S. Ask, J. R. Christiansen, R. Corke, N. Desai, P. Ilten, S. Mrenna, S. Prestel, C. O. Rasmussen, and P. Z. Skands, An introduction to PYTHIA 8.2, *Comput. Phys. Commun.* **191**, 159 (2015);
- [54] T. Han, Z. Si, K. M. Zurek, and M. J. Strassler, Phenomenology of hidden valleys at hadron colliders, *J. High Energy Phys.* **07** (2008) 008.
- [55] L. Carloni and T. Sjöstrand, Visible effects of invisible hidden valley radiation, *J. High Energy Phys.* **09** (2010) 105; L. Carloni, J. Rathsmann, and T. Sjöstrand, Discerning secluded sector gauge structures, *J. High Energy Phys.* **04** (2011) 091.
- [56] S. Mrenna and P. Skands, Automated parton-shower variations in Pythia 8, *Phys. Rev. D* **94**, 074005 (2016); T. Sjöstrand and P. Z. Skands, Transverse-momentum-ordered showers and interleaved multiple interactions, *Eur. Phys. J. C* **39**, 129 (2005).
- [57] J. de Favereau, C. Delaere, P. Demin, A. Giammanco, V. Lemaître, A. Mertens, and M. Selvaggi (DELPHES 3 Collaboration), DELPHES 3, A modular framework for fast simulation of a generic collider experiment, *J. High Energy Phys.* **02** (2014) 057.
- [58] N. Arkani-Hamed and N. Weiner, LHC signals for a superunified theory of dark matter, *J. High Energy Phys.* **12** (2008) 104; M. Baumgart, C. Cheung, J. T. Ruderman, L. T. Wang, and I. Yavin, Non-Abelian dark sectors and their collider signatures, *J. High Energy Phys.* **04** (2009) 014.
- [59] V. Khachatryan *et al.* (CMS Collaboration), Performance of electron reconstruction and selection with the CMS detector in proton-proton collisions at $\sqrt{s} = 8$ TeV, *J. Instrum.* **10**, P06005 (2015); G. Aad *et al.* (ATLAS Collaboration), Electron reconstruction and identification efficiency measurements with the ATLAS detector using the 2011 LHC proton-proton collision data, *Eur. Phys. J. C* **74**, 2941 (2014).
- [60] G. Aad *et al.* (ATLAS Collaboration), Muon reconstruction performance of the ATLAS detector in proton-proton collision data at $\sqrt{s} = 13$ TeV, *Eur. Phys. J. C* **76**, 292 (2016).
- [61] M. Aaboud *et al.* (ATLAS Collaboration), Search for new phenomena in dijet events using 37 fb^{-1} of pp collision data collected at $\sqrt{s} = 13$ TeV with the ATLAS detector, *Phys. Rev. D* **96**, 052004 (2017).
- [62] A. M. Sirunyan *et al.* (CMS Collaboration), Search for narrow and broad dijet resonances in proton-proton collisions at $\sqrt{s} = 13$ TeV and constraints on dark matter mediators and other new particles, [arXiv:1806.00843](https://arxiv.org/abs/1806.00843).
- [63] M. Chala, F. Kahlhoefer, M. McCullough, G. Nardini, and K. Schmidt-Hoberg, Constraining dark sectors with mono-jets and dijets, *J. High Energy Phys.* **07** (2015) 089.
- [64] B. A. Dobrescu and F. Yu, Coupling-mass mapping of dijet peak searches, *Phys. Rev. D* **88**, 035021 (2013); Erratum, **90**, 079901(E) (2014).
- [65] G. Aad *et al.* (ATLAS Collaboration), A search for prompt lepton-jets in pp collisions at $\sqrt{s} = 8$ TeV with the ATLAS detector, *J. High Energy Phys.* **02** (2016) 062.
- [66] G. Aad *et al.* (ATLAS Collaboration), A search for prompt lepton-jets in pp collisions at $\sqrt{s} = 7$ TeV with the ATLAS detector, *Phys. Lett. B* **719**, 299 (2013).
- [67] Y. L. Dokshitzer, G. D. Leder, S. Moretti, and B. R. Webber, Better jet clustering algorithms, *J. High Energy Phys.* **08** (1997) 001; M. Wobisch and T. Wengler, Hadronization corrections to jet cross-sections in deep inelastic scattering, Hamburg 1998/1999, Monte Carlo generators for HERA physics, [arXiv:hep-ph/9907280](https://arxiv.org/abs/hep-ph/9907280).
- [68] M. Cacciari, G. P. Salam, and G. Soyez, FastJet user manual, *Eur. Phys. J. C* **72**, 1896 (2012).
- [69] CMS Collaboration, A search for beyond standard model light bosons decaying into muon pairs, Report No. CMS-PAS-HIG-16-035.